



$t=?$ prehod med prevlado sevanja in prevlado snovi

CMB = cosmic microwave background \rightarrow PRASEVANJE $T_{\text{CMB}} = 2,73\text{K}$

$$\rho_{\text{rad},0} c^2 = 4,2 \cdot 10^{-14} \frac{\text{J}}{\text{m}^3}$$

$$\rho_{\text{m},0} c^2 = 2,6 \cdot 10^{-10} \frac{\text{J}}{\text{m}^3}$$

$$\rho_{\nu} = 0,68 \rho_{\text{rad},0}$$

$$\rho_{\text{m}}(t) c^2 = \rho_{\text{m},0}(t_0) c^2 \frac{R(t_0)^3}{R(t)^3}$$

$$\rho_{\text{rad}}(t) c^2 = \rho_{\text{rad},0} c^2 \frac{R(t_0)^4}{R(t)^4}$$

$$\frac{\rho_{\text{m}}(t)}{\rho_{\text{rad}}(t)} = 1 \rightarrow \frac{\rho_{\text{m},0} c^2}{1,7 \rho_{\text{rad},0} c^2} = \frac{R(t_0)}{R(t)} = 3500 = \frac{t_0^{2/3}}{t^{2/3}}$$

1 (sev) + 0,7 (neutrini)

$$R(t) = \frac{R(t_0)}{3500}$$

Ocenimo starost vesolja

$$t = t_0 (3500)^{-3/2} = \frac{t_0}{200000}$$

Spodnja meja:

• $k=0$ (ravno vesolje)

• kritična gostota $\rho_{0,\text{cr}} = \frac{3H_0^2}{8\pi G}$

• analitična rešitev za prevlado snovi ($\rho \propto R^{-3}$)

$$1.\text{F.E.} \quad \left(\frac{\dot{R}}{R}\right)^2 = \frac{8\pi G}{3} \rho - \frac{k_0 c^2}{R^2} \stackrel{k=0}{=} \frac{8\pi G}{3} \frac{H_0^2}{H_0^2} \rho = H_0^2 \frac{\rho}{\rho_{0,\text{cr}}} = H_0^2 \left(\frac{R_0}{R}\right)^3 = H^2$$

$$\left(\frac{\dot{R}}{R}\right)^2 = H_0^2 \left(\frac{R_0}{R}\right)^3 \rightarrow \frac{\dot{R}}{R} = H_0 \left(\frac{R_0}{R}\right)^{3/2}$$

$$\frac{\dot{R}}{R} = H_0 \left(\frac{R_0}{R}\right)^{3/2} \quad | \cdot R$$

$$\frac{dR}{dt} = H_0 R_0^{3/2} R^{-1/2}$$

$$\int_0^{t_0} dt = \frac{1}{H_0 R_0^{3/2}} \int_0^{R_0} R^{1/2} dR$$

$$\underline{t_0} = \frac{1}{H_0 R_0^{3/2}} \frac{2}{3} R_0^{3/2} = \underline{\frac{2}{3H_0}}$$

Spodnja
meja

Zgornja meja za starost vesolja:

• $p = 0$ (prazno vesolje)

2.F.E.:

$$\frac{\ddot{R}}{R} = -\frac{4}{3} \frac{\pi G}{c^2} (\rho c^2 + 3p)$$

2.F.E. $\ddot{R} = 0$ torej $\dot{R} = \text{konst} = H \cdot R = H_0 R_0$

$$\frac{dR}{dt} = H_0 R_0$$

$$\int_0^{t_0} dt = \frac{1}{H_0 R_0} \int_0^{R_0} dR$$

$$\rightarrow \underline{t_0} = \frac{1}{H_0} \quad \text{Zgornja
meja}$$

Pokazali smo, da za $0 < \Omega_{m,0} < 1$ je starost vesolja

$$\frac{2}{3H_0} < t_0 < \frac{1}{H_0}$$

$$(H_0 = 70 \text{ km/s/Mpc}) \quad 9 \text{ Gyr} < t_0 < 14 \text{ Gyr}$$

$$t = \frac{t_0}{200000} \sim 65000 \text{ let}$$

FRIEDMANOVE ENAČBE S KOZMOLOŠKO KONSTANTO

$$\left(\frac{\dot{R}}{R}\right)^2 = \frac{8\pi G}{3} \rho - \frac{kc^2}{R^2} + \frac{\Lambda c^2}{3} \quad 1.F.E.$$

$$\frac{\ddot{R}}{R} = -\frac{4\pi G}{3c^2} (\rho c^2 + 3p) + \frac{\Lambda c^2}{3} \quad 2.F.E.$$

$$\dot{\rho} c^2 = -3 \frac{\dot{R}}{R} (\rho c^2 + p) \quad 3.F.E.$$

Einstein: statična
rešitev: $\dot{R} = \ddot{R} = 0$

$$\vee 1.F.E \quad \frac{8\pi G}{3} \rho + \frac{\Lambda c^2}{3} = \frac{8\pi G}{3c^2} \left(\rho c^2 + \frac{\Lambda c^4}{8\pi G} \right)$$

Kaj se dogaja do zelo velikih časih?

$$\left(\frac{\dot{R}}{R}\right)^2 = \frac{8\pi G}{3} \rho - \frac{kc^2}{R^2} + \frac{\Lambda c^2}{3} \quad \Rightarrow \quad H^2 = \left(\frac{\dot{R}}{R}\right)^2 \sim \frac{\Lambda c^2}{3}$$

$\rho \propto \dot{R}^{-3}$ $\frac{kc^2}{R^2} \propto R^{-2}$

$$\dot{R} \sim \sqrt{\frac{\Lambda c^2}{3}} R$$

$$\frac{dR}{R} \sim \sqrt{\frac{\Lambda c^2}{3}} dt \quad \rightarrow \quad R(t) \propto e^{\sqrt{\frac{\Lambda c^2}{3}} t} = e^{Ht}$$

REŠITEV ZA ŠIRJENJE VESOLJA S TEMNO ENERGIJO

Opazovanja: $\Lambda > 0$ (konst. vrednost)

$k = 0$

$$\left(\frac{\dot{R}}{R}\right)^2 = \frac{8\pi G}{3} \rho - \frac{k c^2}{R^2} + \frac{\Lambda c^2}{3}$$

$$H^2 = \frac{8\pi G}{3} \rho + \frac{\Lambda c^2}{3} \quad | : H_0^2$$

$$\frac{H^2}{H_0^2} = \frac{8\pi G}{3H_0^2} \rho + \frac{\Lambda c^2}{3H_0^2}$$

$$\frac{H^2}{H_0^2} = \frac{\rho}{\rho_{cr,0}} + \frac{\Lambda c^2}{3H_0^2}$$

Ω_Λ gostotni parameter za kozmološko konst.

meritev

$$\rho_{cr,0} = \frac{3H_0^2}{8\pi G}$$

opazovanja!

parameter gostote $\Omega_m = \frac{\rho_m}{\rho_{cr}}$
 $m = \text{mater}$

$$\frac{H^2}{H_0^2} = \Omega_m + \Omega_\Lambda$$

ob $t = t_0 \rightarrow \Omega_{m,0} + \Omega_{\Lambda,0} = 1$

- če je $k=0$ (ravno vesolje) velja ob kateremkoli času $\Omega_m + \Omega_\Lambda = 1!$
- če je $k=0$ in $\Lambda=0 \quad \Omega_m = 1!$
- če je $\Lambda > 0$ in k poljuben ($k \neq 0$):

$\Omega_m + \Omega_\Lambda > 1$ zaprto vesolje
 $\Omega_m + \Omega_\Lambda < 1$ odprto vesolje

pregledam

število vesolja

$$\Omega_m = \frac{\rho_{m,0}}{\rho_{cr,0}} = 1$$

H_0 rdeči premik galaktij

$$R \propto \frac{1}{1+z} \quad \left(\text{glej profinjico: } z = \frac{R_0}{R} - 1 \rightarrow \frac{R}{R_0} = \frac{1}{1+z} \right)$$

$$R(t_0) \propto \frac{1}{1+0}$$

$$R(t_1) \propto \frac{1}{1+z(t_1)}$$

$$\frac{R}{R_0} = \frac{1}{1+z}$$

$$R(t_2) \propto \frac{1}{1+z(t_2)}$$

$$\left. \begin{array}{l} R(t_1) \\ R(t_2) \end{array} \right\} = \frac{1+z(t_2)}{1+z(t_1)}$$

1. F. E. z $\frac{R}{R_0} = \frac{1}{1+z}$ lahko zapišemo kot

$$H = \frac{\dot{R}}{R} = \frac{\dot{a}}{a} =$$

$$a \rightarrow 0 \quad z \rightarrow \infty$$

$$a = 1 \quad z = 0$$

skalirni faktor

$$= H_0 \sqrt{(\Omega_b + \Omega_c) \bar{a}^3 + \Omega_{rad} \bar{a}^4 + \Omega_k \bar{a}^2 + \Omega_\Lambda \bar{a}^{-3(w+1)}}$$

barionska
snov

hladna
temna snov

$$\Omega_m = \frac{\rho}{\rho_{cr}}$$

$$p_\Lambda = w p_\Lambda c^2$$

$w = -1$
za kozm.
konst.

Rešitev v primeru:

$k=0, \Lambda > 0 (w=-1), \Omega_{rad}$ zanemarimo:

$$H(a) = \frac{\dot{a}}{a} = \sqrt{\Omega_m \bar{a}^3 + \Omega_\Lambda}$$

$w \neq -1$
krivina

$$a(t) = \left(\frac{\Omega_m}{\Omega_\Lambda} \right)^{1/3} \sinh^{2/3} \left(\frac{t}{t_\Lambda} \right) \quad ; \quad t_\Lambda = \frac{2}{3H_0 \sqrt{\Omega_\Lambda}}$$

(uporabna za $a > 0,01$ ($t > 10^7$ let))

od kdaj se tako vesolje širi pospešeno?

$$a(\ddot{a}=0) = \left(\frac{\Omega_m}{2\Omega_\Lambda} \right)^{1/3}$$

$$a \sim 0,6$$

$$z \sim 0,66$$

OPAZOVALNI PREIZKUSI

Glavna opazovanja, ki določajo kozmološke parametre, so:

- obstoj in struktura prahuvanja
- porazdelitev snovi na velikih skalah
- kemijska zastopanost vodika, helija in litija ← prvinska nukleosinteza
- pospešeno širjenje vesolja z oddaljenimi supernovami Ia

Kozmološki rdeči premik

FRW metrika: za fotone $ds = 0$

$$0 = c^2 dt^2 - R(t)^2 \frac{dr^2}{1 - kr^2} \quad (\theta = \phi = 0)$$

① r_e

$\rightsquigarrow t_e$

$\rightsquigarrow t_e + \Delta t_e$

$r = 0$ ②

t_0

$t_0 + \Delta t_0$

prvi foton

$$\int_{t_e + \Delta t_e}^{t_0 + \Delta t_0} \frac{dt}{R(t)} = \frac{1}{c} \int_0^{r_e} \frac{dr}{\sqrt{1 - kr^2}}$$

drugi foton

$$\int_{t_e + \Delta t_e}^{t_0 + \Delta t_0} \frac{dt}{R(t)} - \int_{t_e}^{t_0} \frac{dt}{R(t)} = 0$$

$$\int_{t_e}^{t_0} - \int_{t_e}^{t_e + \Delta t_e} + \int_{t_0}^{t_0 + \Delta t_0} - \int_{t_e}^{t_0} \frac{dt}{R(t)} = 0$$

$$\frac{\Delta t_e}{R(t_e)} = \frac{\Delta t_0}{R(t_0)}$$

$$\frac{\Delta t_e}{\Delta t_0} = \frac{\lambda_e}{\lambda_0} = \frac{c}{c}$$

$$\frac{\Delta t_0}{\Delta t_e} = \frac{\lambda_0}{\lambda_e} = \frac{R(t_0)}{R(t_e)} = 1 + z$$

Razdalje:

- gostota svetlobnega toka
razdalja izsera d_L

$$j = \frac{L}{4\pi d_L^2} = \frac{L}{4\pi d^2 (1+z)^2}$$

$$d_L = d(1+z) \underset{k=0}{=} a \cdot r(1+z)$$

- sprememba en. fotonov
- fotoni prihajajo manj pogosto

- kotna velikost objekta

razdalja kotnega premera d_{premu}

$$l = r_0 a(t_e) \theta$$

$$\theta = \frac{l}{r_0 a(t_e)} = \frac{l(1+z)}{r_0 a_0}$$

$$d_{\text{premu}} = \frac{l}{\sin \theta} \sim \frac{l}{\theta} = \frac{r_0 a_0}{1+z} = \frac{d_L}{(1+z)^2}$$

$$d_{\text{premu}} = \frac{d}{1+z} = \frac{d_L}{(1+z)^2}$$

Mikrovalotno seraije otadja

$$\left. \begin{aligned} p_{\text{rad}} &\propto \frac{1}{a^4} \\ p_{\text{rad}} &= \left(\frac{4}{c}\right) \sigma T^4 \end{aligned} \right\} T \propto \frac{1}{a}$$

$$T(t_0) = 2,73 \text{ K}$$

$$T \sim 3000 \text{ K}$$